

# Moment methods and nuclear level densities

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supported by a SSAA grant from US DOE-NNSA

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## Tools: mean-fields and moments

3 approaches to nuclear level density:

modified Fermi gas

Monte Carlo shell model

spectral distribution methods



mean-field (centroids or first moments)

residual interaction (spreading widths or second moments)

collective interaction (third moments)

## Fermi Gas Models

Single-particle energies from Hartree-Fock *mean-field*:  $\varepsilon_i^{p,n}$

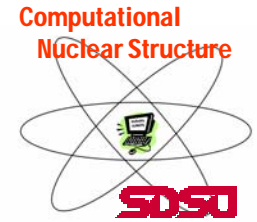
Single-particle density of states:  $g(\varepsilon) = \sum_i \delta(\varepsilon - \varepsilon_i)$

Partition function :  $\ln Z(\alpha, \beta) = \sum_i \ln(1 + \exp(\alpha - \beta\varepsilon_i))$

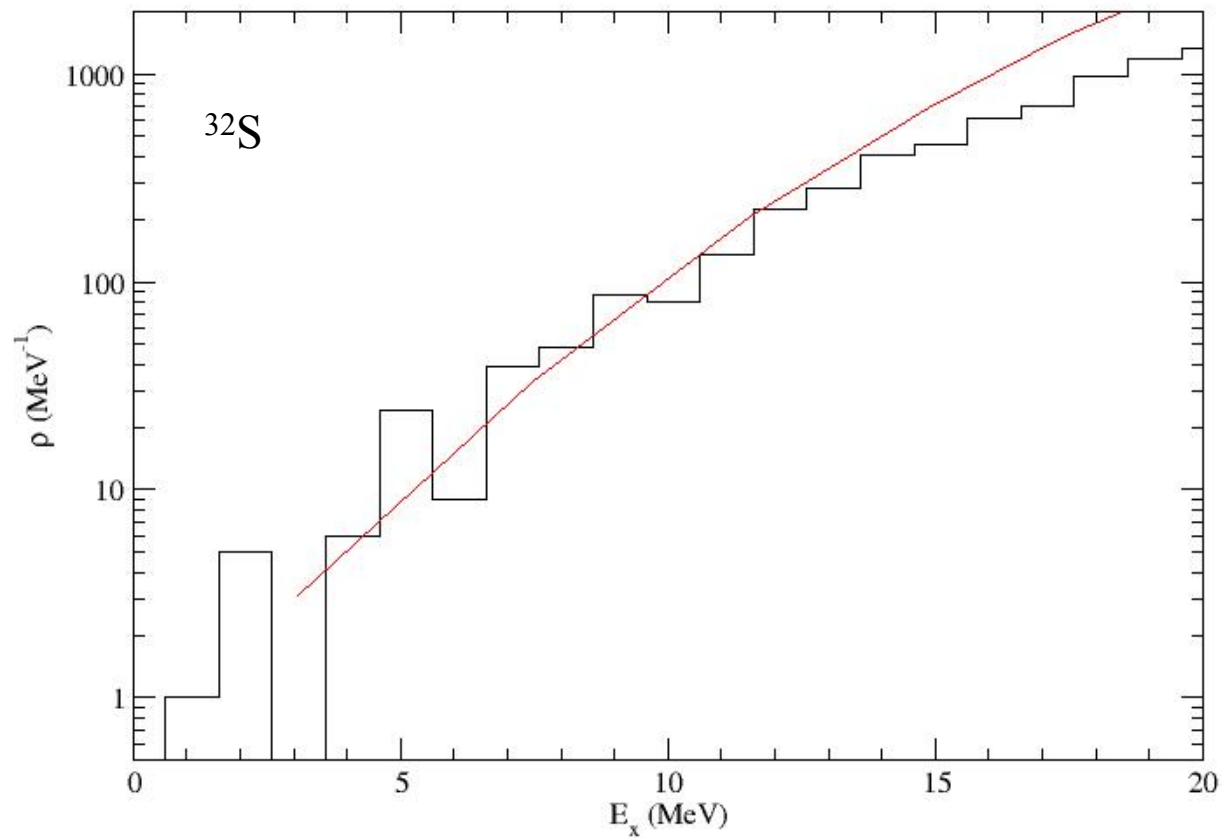
Then apply saddle-point method...

**Please** keep in mind the following is **deliberately crude** and **unsophisticated** and is meant as **motivation**, not **criticism** of this approach

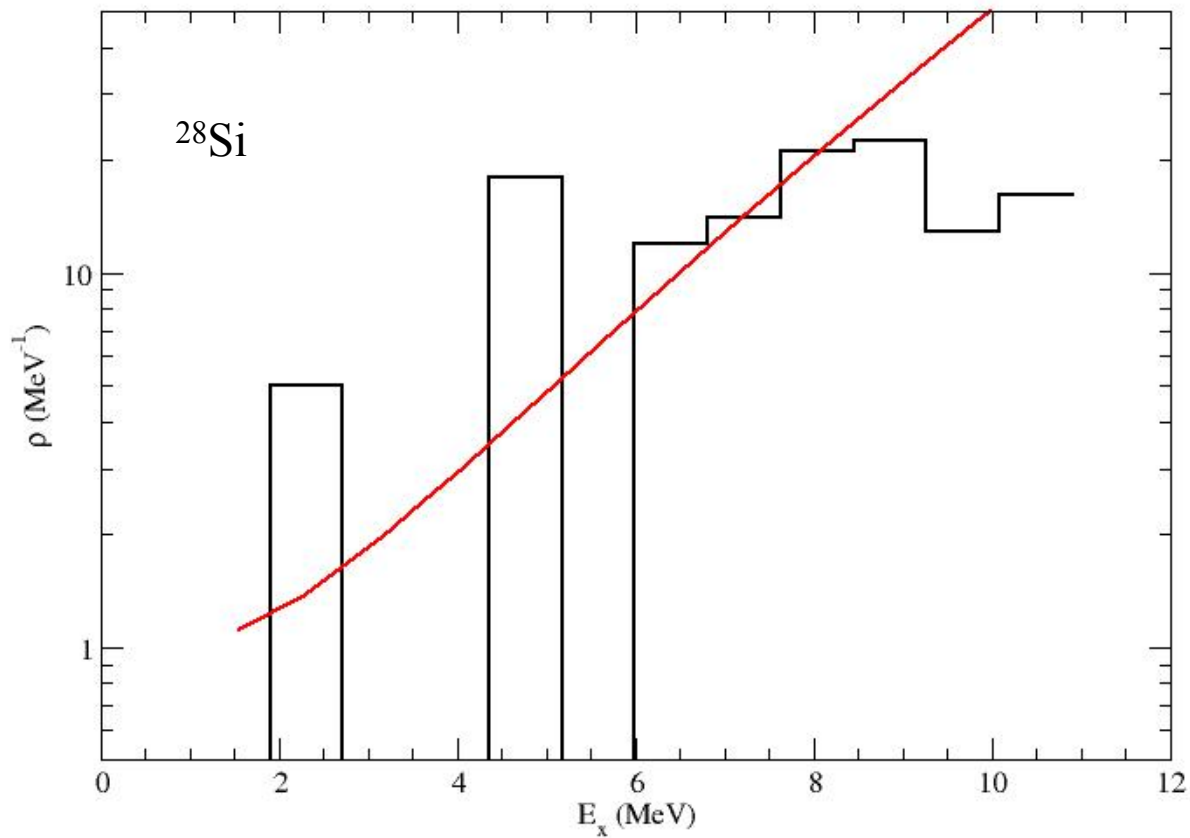
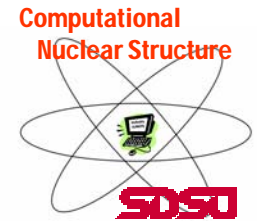
# Mean-field Level densities



First, start with “exact” level density from shell-model diagonalization and add in model of level density from mean-field calculation

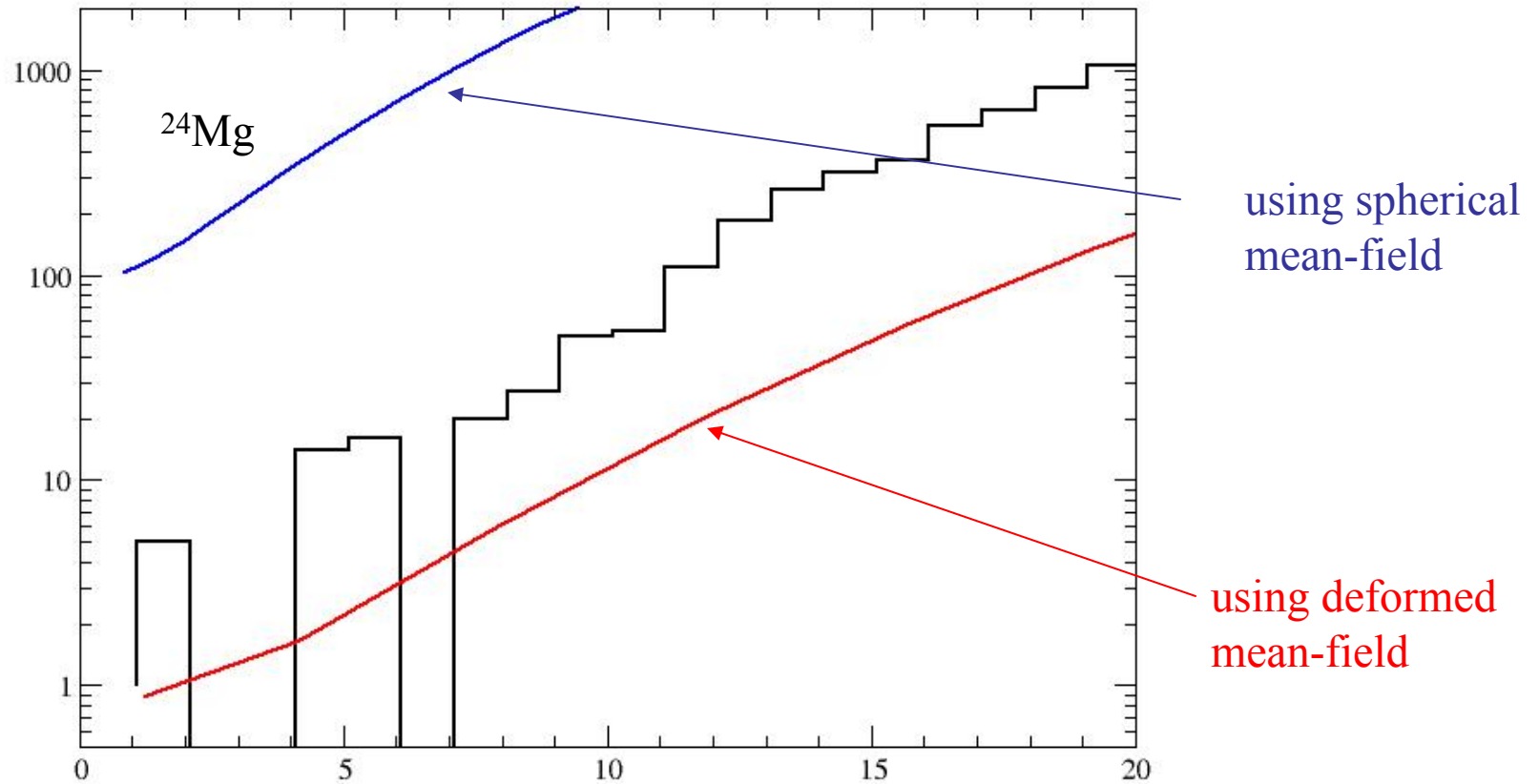
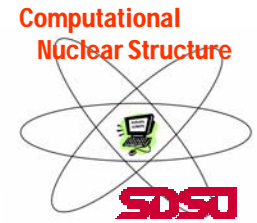


# Mean-field Level densities

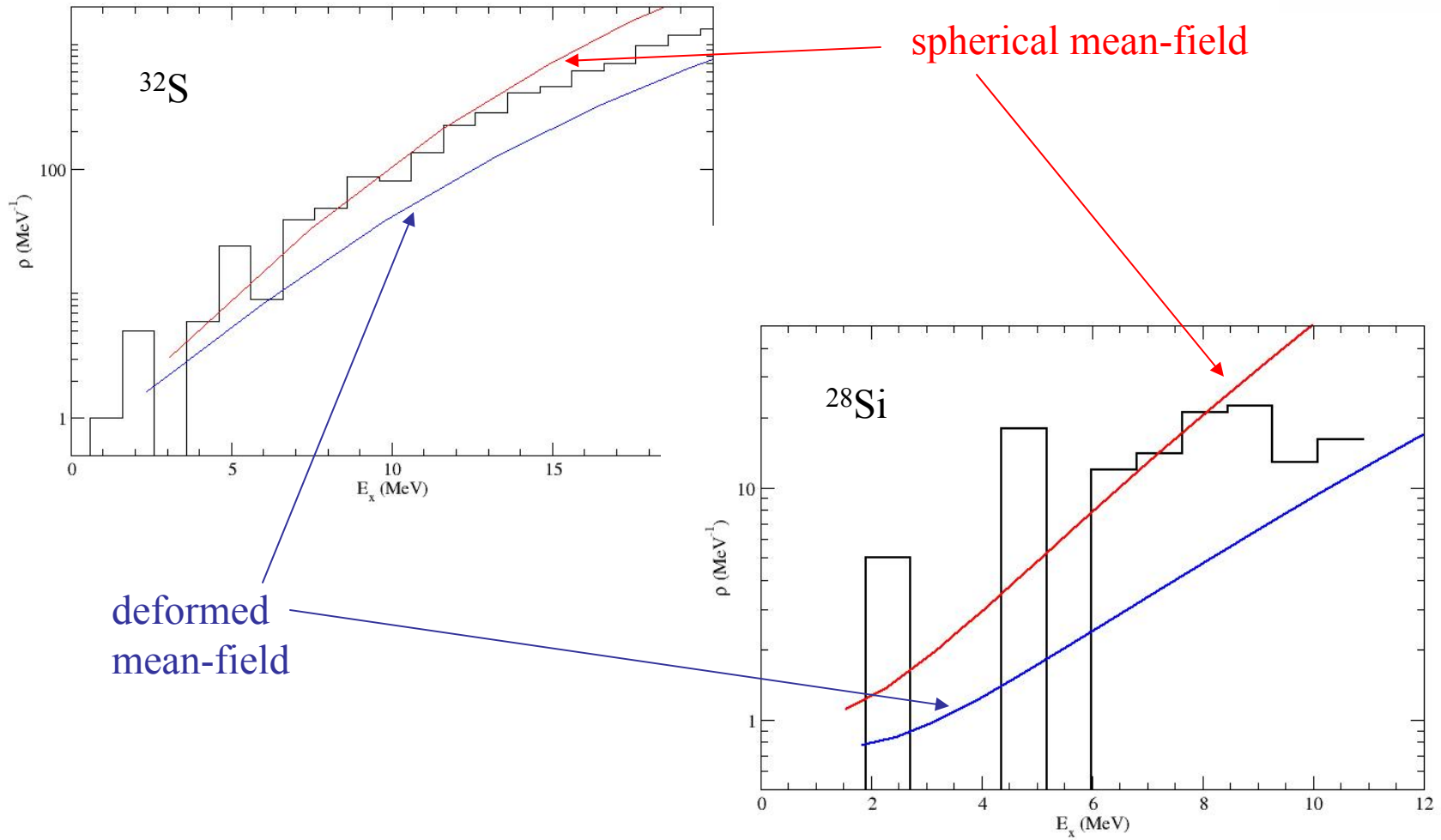


These are both  
spherical nuclei...  
what about  
deformed nuclides?

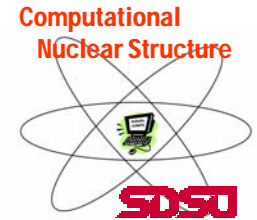
# Mean-field Level densities



# Mean-field Level densities

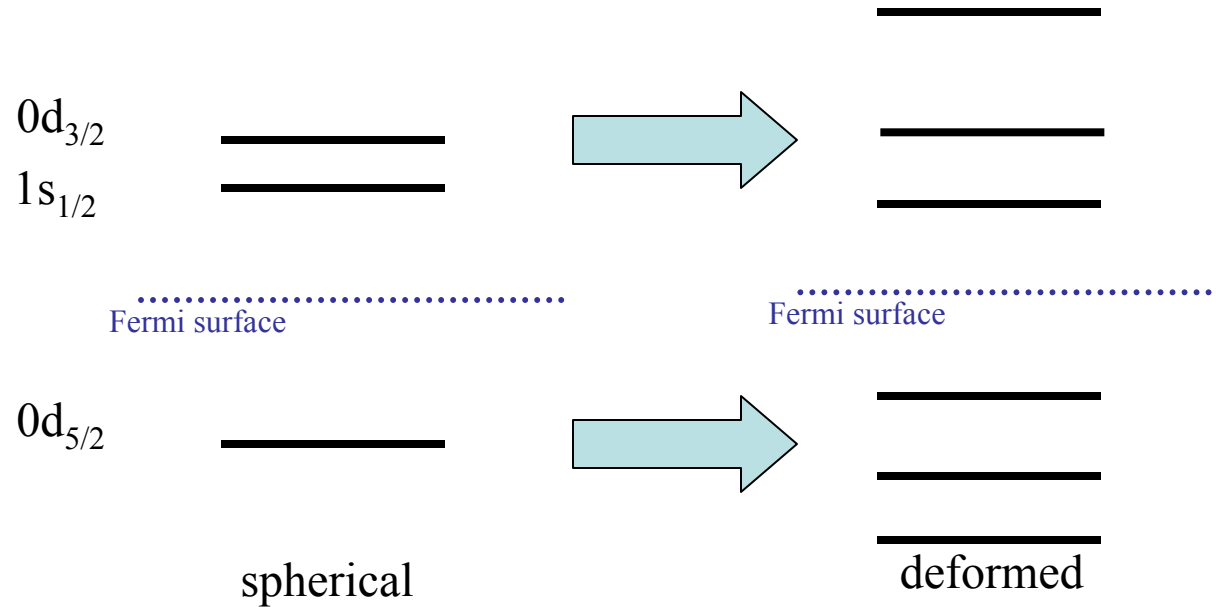


# Mean-field Level densities



Difference is due to **fragmentation** of Hartree-Fock single-particle energies in deformed mean-field

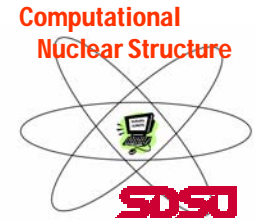
This in turn is a manifestation of the **residual interaction**





# Introduction to Statistical Spectroscopy

(also known as “spectral distribution theory”)



Pioneered by J. Bruce French 1960's-1980's

*other luminaries include: J. P. Draayer, J. Ginocchio, S. Grimes, V. Kota, S.S.M. Wong, A.P. Zuker + many others...*

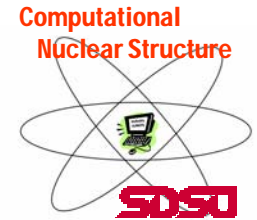
Problem: diagonalization is too hard and gives too much detailed information

Solution: instead of diagonalizing  $\mathbf{H}$ , find moments:  $\text{tr } \mathbf{H}^n$

Key question: how many moments do we need?

Rather than many moments (over the entire space)  $\text{tr } \mathbf{H}^n$ ,  $n = 1, 2, 3, 4, 5, 6, 7, \dots$   
compute low moments ( $n = 1, 2, 3, 4$ ) on *subspaces*

# How we do it: a detailed version



## The important configuration moments

Dimension  $d_\alpha = \text{Tr} \mathbf{P}_\alpha$

Centroid:  $\bar{E}_\alpha = \frac{1}{d_\alpha} \text{Tr} \mathbf{P}_\alpha \mathbf{H}$       Width:  $\sigma_\alpha = \frac{1}{d_\alpha} \text{Tr} \mathbf{P}_\alpha (\mathbf{H} - \bar{E}_\alpha)^2$

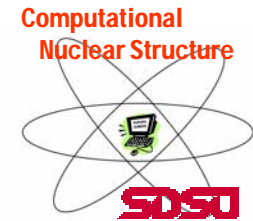
Higher central moments  $\mu_n(\alpha) = \frac{1}{d_\alpha} \text{Tr} \mathbf{P}_\alpha (\mathbf{H} - \bar{E}_\alpha)^n$

Scaled moments  $m_n(\alpha) = \mu_n(\alpha) / (\sigma_\alpha)^n$

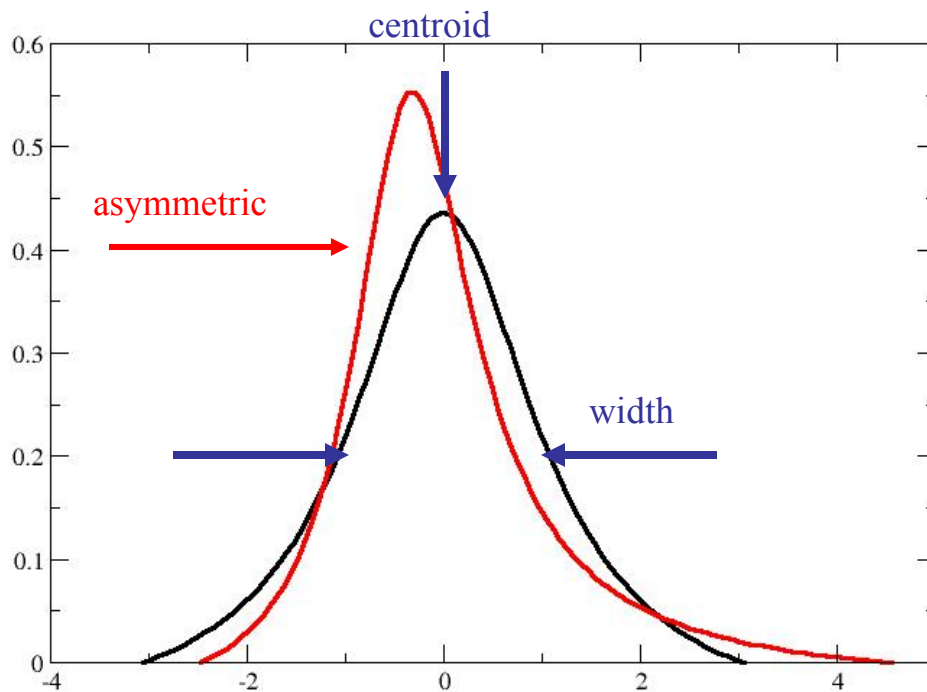
Asymmetry (or skewness):  $m_3(\alpha)$

Excess :  $m_4(\alpha) - 3 = 0$  for Gaussian

# Introduction to Statistical Spectroscopy



## Primer on moments



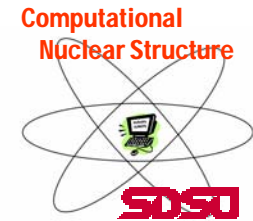
### **Interpretation** of moments:

centroid = spherical HF energy

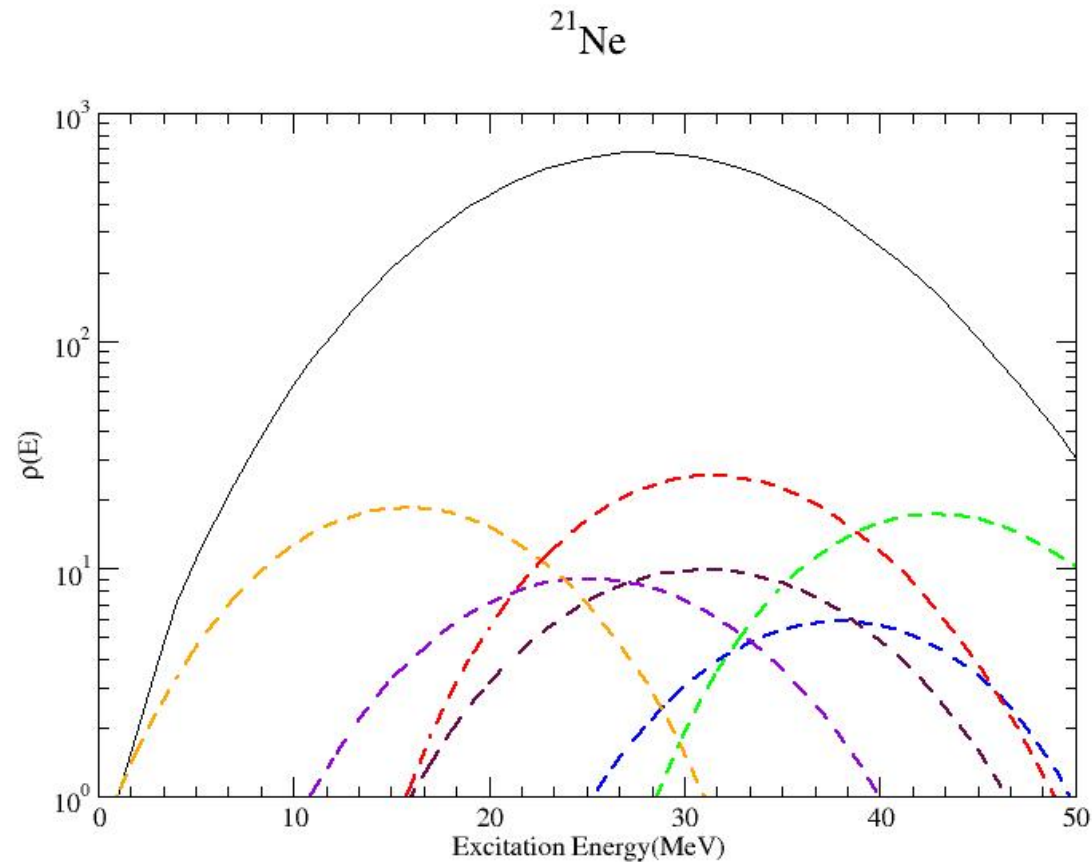
width = avg spreading width  
of residual interaction

asymmetry = measure of collectivity

# Introduction to Statistical Spectroscopy

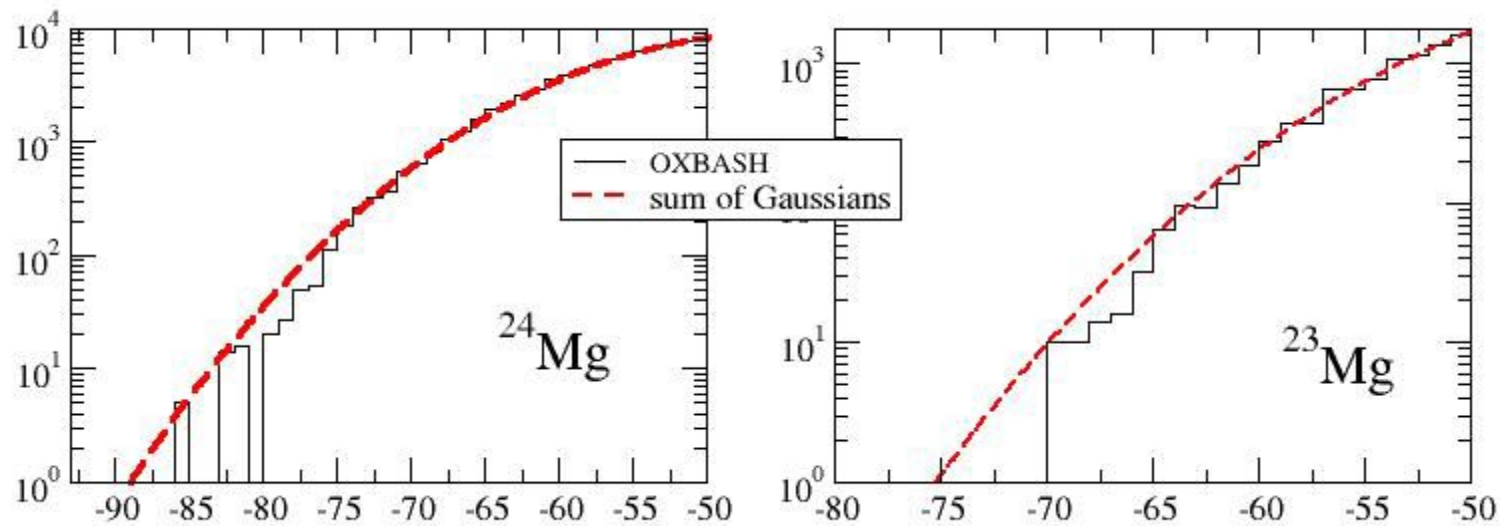


Then we consider the level density as being the sum of individual configuration densities

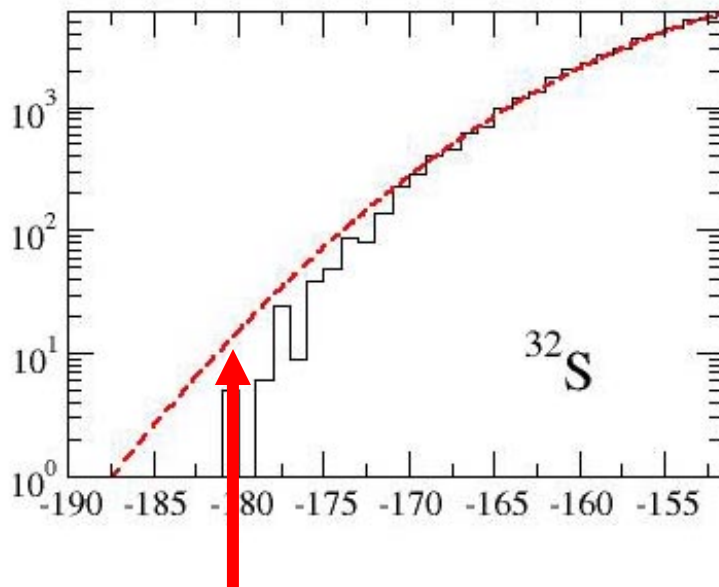


## Level densities as a sum of configuration densities

We model the level density as a sum of partial (configuration) densities, each of which are modeled as Gaussians



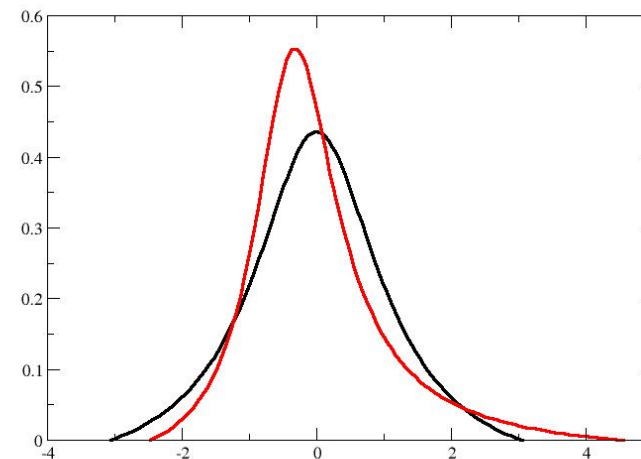
## Level densities as a sum of configuration densities



Not satisfactory!

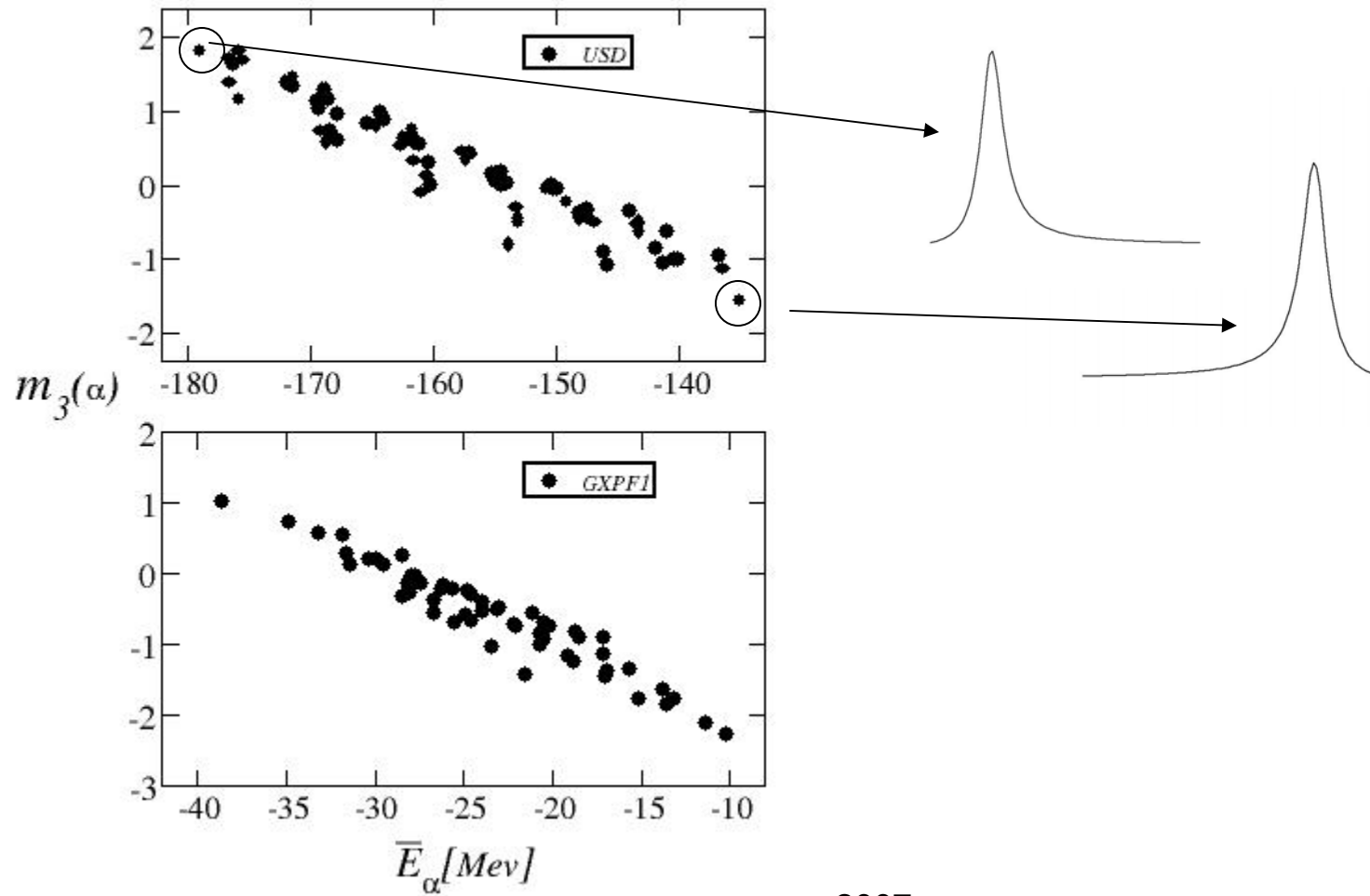
What can we do to  
improve our model?

Go to third moments: asymmetries

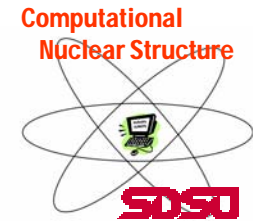


# Shell-Model Configuration moments

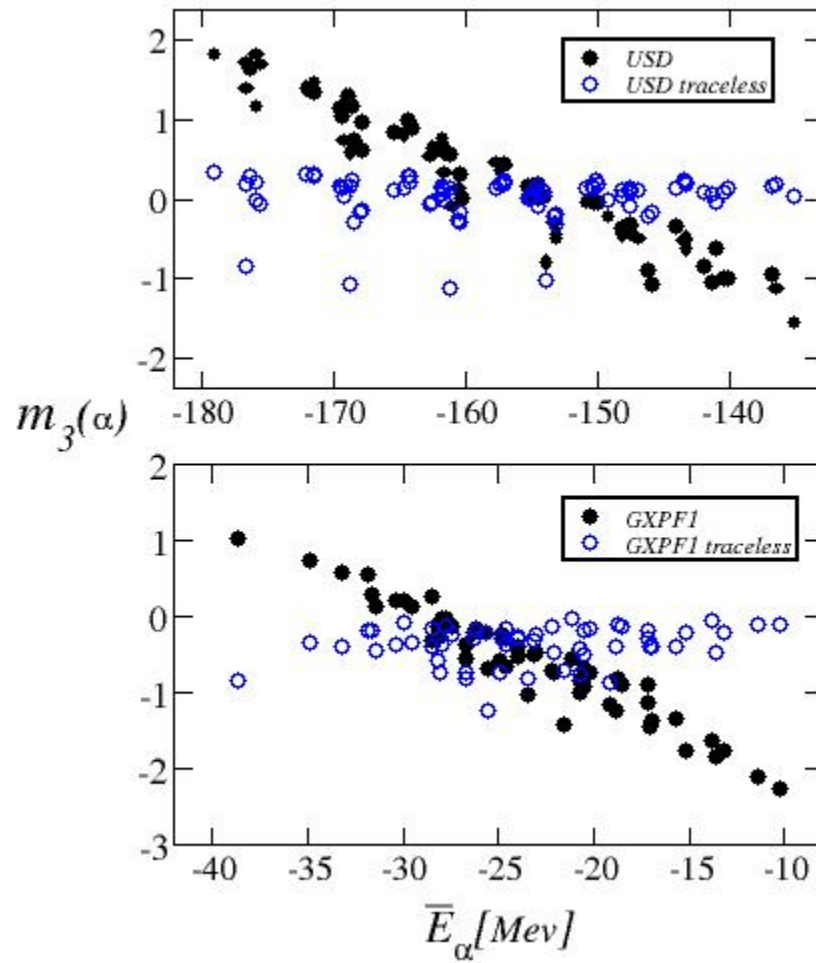
The configuration asymmetry varies almost linearly with the centroid



# Shell-Model Configuration moments



The configuration asymmetry varies almost linearly with the centroid



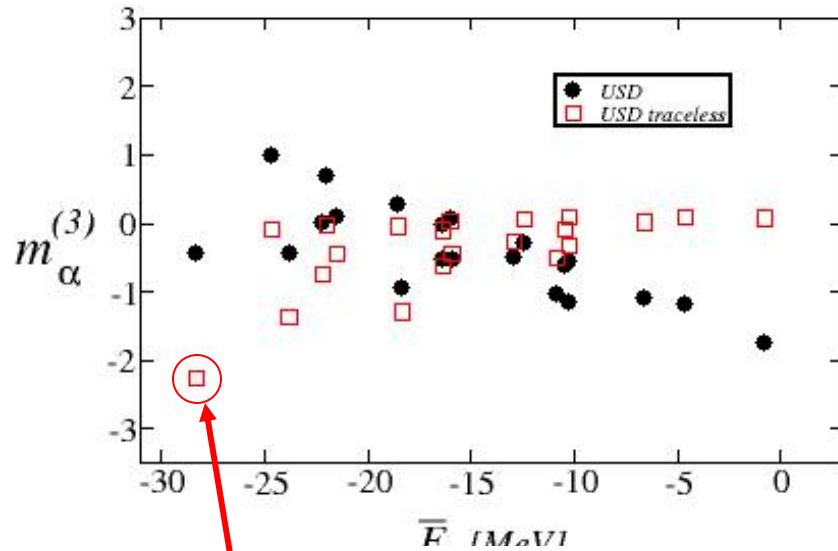
We can make all the centroids =0  
by setting the “monopole” part  
of the interaction = 0  
(this is called a “traceless” interaction in  
the vocabulary of statistical spectroscopy)

The monopole potential is related  
to the mean field!

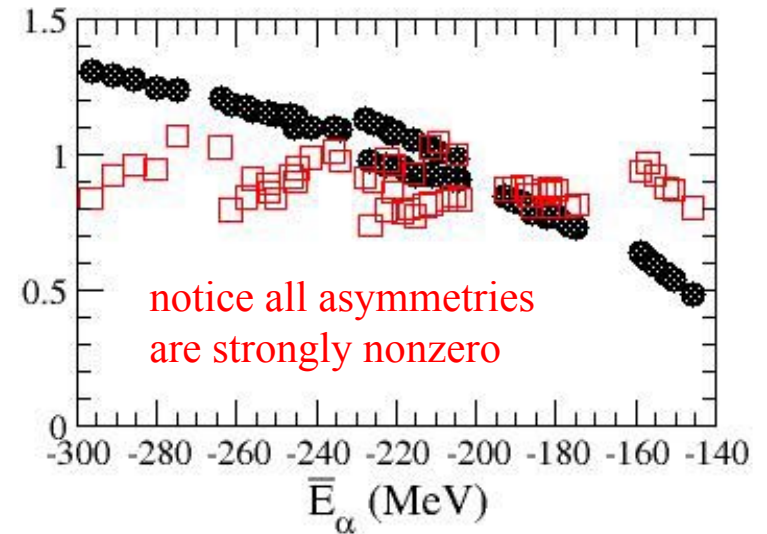
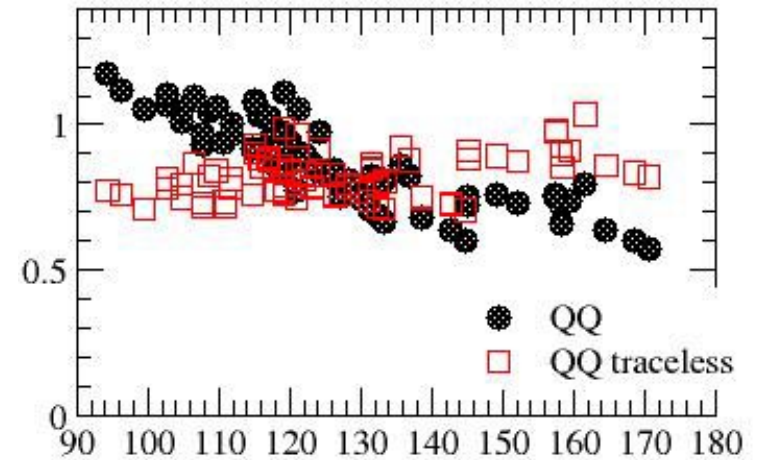
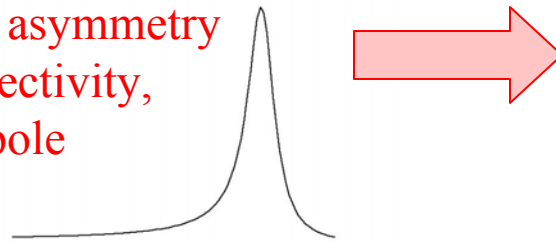
Deviations from this trend are  
associated with strong collectivity



# Collectivity and Asymmetries



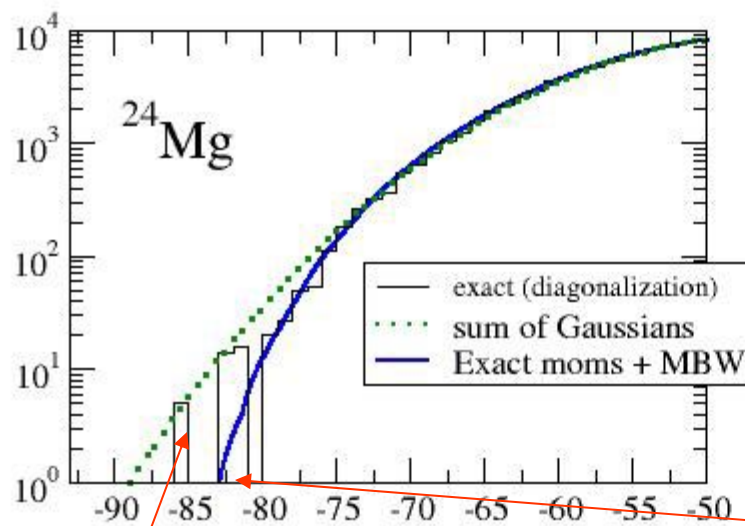
strong traceless asymmetry  
arises from collectivity,  
such as quadrupole



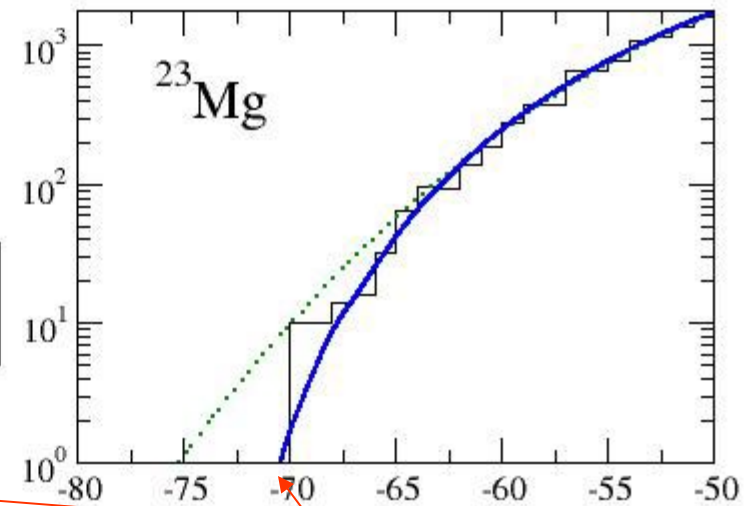
## Level densities as a sum of configuration densities

It is (often) important to include  
3<sup>rd</sup> and 4<sup>th</sup> moments

much better than  
using only second moments

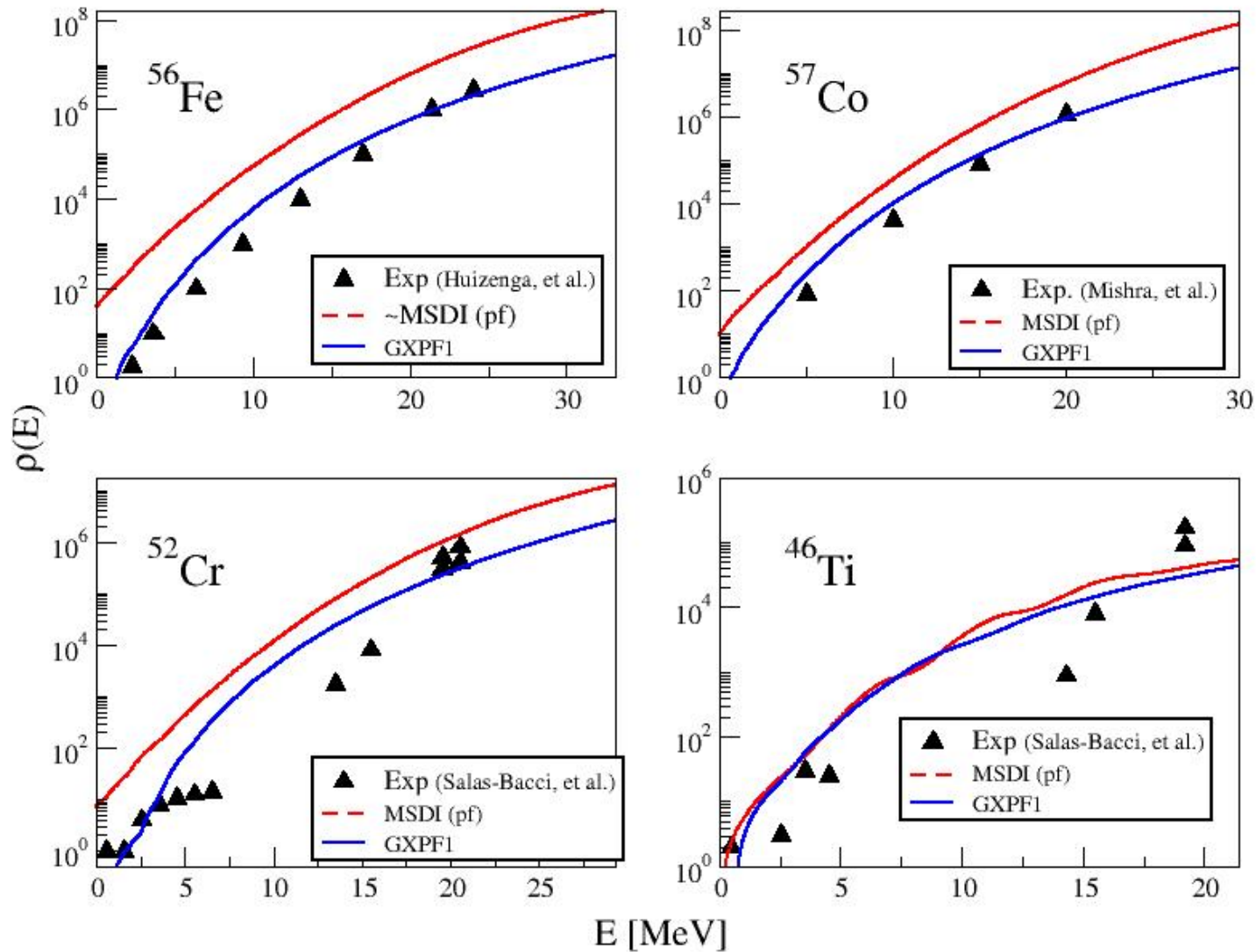


collective states difficult to get



“starting energy” also difficult to control

## Comparison with experiments



NB: computed  
+ parity states  
and multiplied  $\times 2$

## Obligatory Summary

View nuclear many-body Hamiltonian through lens  
of moment methods:

1<sup>st</sup> (configuration) moments = mean-field

2<sup>nd</sup> moments = spreading widths of residual interaction

3<sup>rd</sup> moments = collectivity of residual interaction